

## Dependence of microwave surface impedance on crystallographic orientation in $\text{YBa}_2\text{Cu}_3\text{O}_{7-x}$ thin films

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**Abstract** : We report the variation of microwave surface resistance ( $R_s$ ) and residual penetration depth  $\lambda(0)$  of YBCO thin films as a function of orientation of the films. Three thin film samples A, B and C have been studied. As revealed through XRD patterns, film A is partially  $c$ -axis oriented, film B has better  $c$ -axis orientation while film C is fully  $c$ -axis oriented. The modified end plate replacement technique operating at 10 GHz in  $\text{TE}_{011}$  mode and sapphire loaded dielectric resonator technique operating at 13.6 GHz in  $\text{TE}_{01\delta}$  mode were used for the measurement of  $R_s$  and  $\lambda(0)$ . The value of  $R_s$  (13.6 GHz, 65 K) for samples A, B and C have been found to be 1.607, 1.356 and 1.037 m $\Omega$  respectively. Lower values of  $R_s$  (65 K) = 930, 780 and 600  $\mu\Omega$  have been measured at 10 GHz for these films. The values of  $\lambda(0)$  (13.6 GHz) have been found to be 1950, 1830 and 1650 Å for the films A, B and C respectively. As the degree of orientation of  $c$ -axis grains increases in thin films, the Josephson coupling gets stronger and the value of  $R_s$  and  $\lambda(0)$  decreases.

**Keywords** : Microwave surface resistance, YBCO thin films, grain orientation

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### 1. Introduction

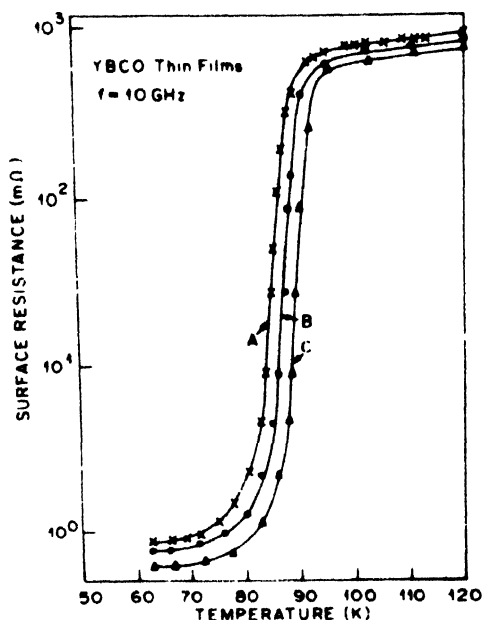
The crystallographic defects in metallic conductors have been a source of power loss from dc to electromagnetic wave frequencies and their control with the increase of purity from 99.5–99.9995% for hyperconductors accompanied by large increase in the resistivity ratio  $\rho$  (300 K)/ $\rho$  (20 K) from 50–6000 is reported in the literature. Low- $T_c$  metallic superconductors with low values of  $R_s$  (10 GHz)  $\approx$  25  $\mu\Omega$  and  $\lambda(0) \approx$  1400 Å at liquid helium temperature have however been in use in microwave superconducting electronics for quite some time. With the discovery of high- $T_c$  cuprate superconductors, various studies have been conducted on YBCO ceramic material processed under different conditions and

whereas in the sample B, the intensity of the 113, 116/123 and 213 peaks diminish and peaks corresponding to the 103/110 and (00n) planes only appear. The film C contains peaks corresponding to (00n) planes only. Film A is partially *c*-axis oriented, film B has better *c*-axis orientation while film C is fully *c*-axis oriented. Other characteristics of the samples are given in the Table 1. The variation of  $R_s$  of the samples with temperature at

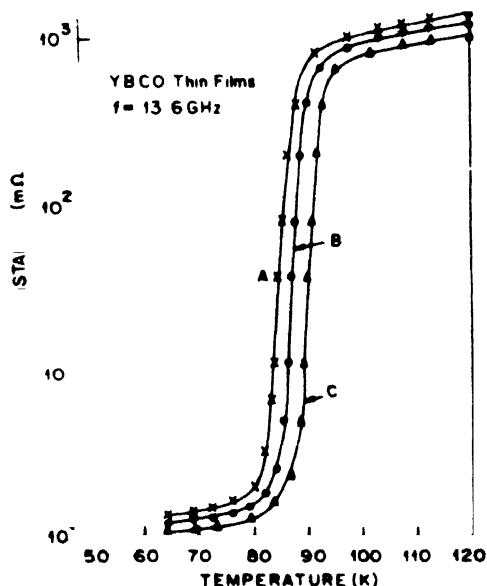
**Table 1.** Characteristics of the samples

Samples	A	B	C
$T_C$ (K)	87	88	90
$\Delta T_C$ (K)	1.7	1.5	1.4
$R_s$ (10 GHz, 65 K) ( $\mu\Omega$ )	930	780	600
$R_s$ (13.6 GHz, 65 K) (m $\Omega$ )	1.607	1.356	1.037
$\lambda(0)$ (13.6 GHz) ( $\text{\AA}$ )	1950	1830	1650

10 GHz is given in Figure 2 whereas the corresponding results at 13.6 GHz are shown in Figure 3. The rate of fall of the microwave surface resistance ( $dR_s/dT$ ) just below  $T_C$  has



**Figure 2.** Variation of  $R_s$  (10 GHz) of the samples with temperature.



**Figure 3.** Variation of  $R_s$  (13.6 GHz) of the samples with temperature.

been found to become steeper as the number of *c*-axis oriented grains increases in the film samples A to B to C and the value of  $R_s$  (65 K) has been found to be the lowest for the highly *c*-axis oriented sample C. The variation of the resonant frequency of the dielectric resonator with temperature is shown in Figure 4. The computed values of  $\lambda(0)$  have been

### Dependence of microwave surface impedance etc

found in the decreasing order in the samples A to B to C. The high- $T_C$  superconductor is an inhomogeneous medium of superconducting grains coupled with Josephson weak links.

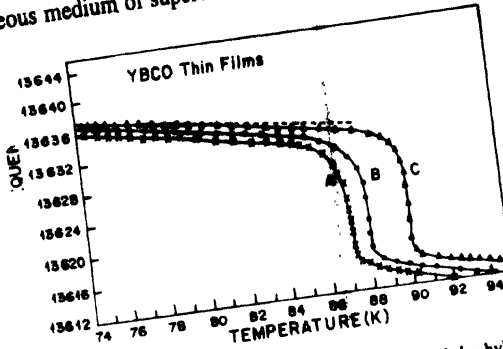


Figure 4. Variation of the resonant frequency ( $f_0$ ) of the hybrid dielectric resonator with temperature.

The total surface resistance of the YBCO thin films in the superconducting state is given by eq. (4)

$$R_s(\text{Total}) = R_s(\text{Intragrain}) + R_s(\text{Intergrain}) + R_s(\text{others}). \quad (4)$$

As the temperature of the YBCO thin film sample is lowered below  $T_C$  (onset), the contribution due to  $R_s(\text{Intragrain})$  and  $R_s(\text{Intergrain})$  decreases and consequently  $R_s(\text{Total})$  decreases.  $R_s(\text{Intergrain})$  is dependent upon the strength of the Josephson coupling. The increase of Josephson coupling strength through increase of orientation of the grains has resulted in the lower value of  $R_s$ . In the randomly oriented samples, the measured value of the penetration depth  $\lambda(T)$  is the function of the penetration depth  $\lambda_L$  in the superconducting grains and  $\lambda_J$  in the weak links. The value of  $\lambda_J$  is also dependent upon the strength of the Josephson coupling. As the degree of the orientation of the  $c$ -axis grains improves, the value of  $\lambda_J$  decreases. The value of  $\lambda(0)$  has been found in the samples A, B and C equal to 1950, 1830 and 1650 Å respectively.

#### 4. Conclusions

The study of  $R_s$  (10 and 13.6 GHz) and  $\lambda(0)$  (13.6 GHz) conducted on oriented thin film samples suggests that with the grain alignment, the weak links are eliminated to a significant level and Josephson coupling gets tightened and the reduced values of these parameters are observed. The minimum achieved value of  $R_s$  (10 GHz, 65 K) for fully  $c$ -axis oriented thin film sample C is 600  $\mu\Omega$ , which could be reduced further reaching residual  $R_s(0)$  value at 0 K. The value of residual penetration depth  $\lambda(0) = 1650$  Å achieved for sample C is however, close to the London's microscopic value of 1400 Å.

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## References

- [1] G Muller, D J Brauer, R Eujen, M Hein, N Klein, H Piel, L Ponto, U Klein and M Peiniger *IEEE Trans. Magn.* **25** 2402 (1989)
- [2] Q Li, K W Rigby and M S Rzechowski *Phys. Rev.* **B39** 6607 (1989)
- [3] A M Awasthi, J P Carini, B Alavi and G Gruner *Solid State Commun.* **67** 373 (1988)
- [4] J Wosik, R A Kranenburg, J C Wolfe, V Selvamanickam and K Salama *J. Appl. Phys.* **69** 874 (1991)
- [5] Ram Swarup and A K Gupta *Indian J. Pure Appl. Phys.* **32** 464 (1994)
- [6] A K Gupta, Ram Swarup and M C Bansal *Physica C* **235-240** 2005 (1994)
- [7] Ram Swarup, A K Gupta and M C Bansal *J. Supercond.* **8** 361 (1995)
- [8] N Klein, G Muller, S Orbach, H Piel, H Chaloupka, B Roas, L Schultz, U Klein and M Peiniger *Physica C* **162-164** 1549 (1989)
- [9] R Pinto, S P Pai, C P D'souza, L C Gupta, R Vijayaraghavan, D Kumar and M Sharon *Physica C* **196** 264 (1992)
- [10] R Chandra, A K Gupta and V Kumar *Indian J. Pure Appl. Phys.* **32** 133 (1994)
- [11] C Le Paven-Thivet, M Guilloux-Viry, J Padiou, A Perrin, M Sergent, L A de Vaulchier and N Bontemps *Physica C* **244** 231 (1995)
- [12] X Y Li, N J Wu, K Xie, J S Liu, H Lin, T Q Huang and A Ignatiev *Physica C* **248** 281 (1995)
- [13] Ram Swarup, A K Gupta, S K Srivastava, V Kumar and M C Bansal *Solid State Phys. (India)* **38C** 284 (1995)
- [14] Ram Swarup, A K Gupta and M C Bansal *Indian J. Cryog.* **20** 3 (1995)